

Hamiltonian Decoded Quantum Interferometry for General Pauli Hamiltonians

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Math Picture Language Seminar @ Harvard

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Based on joint work with Weichen Gu, Dax Koh, Xiang Li
[arXiv:2508.10725, 2601.18773]

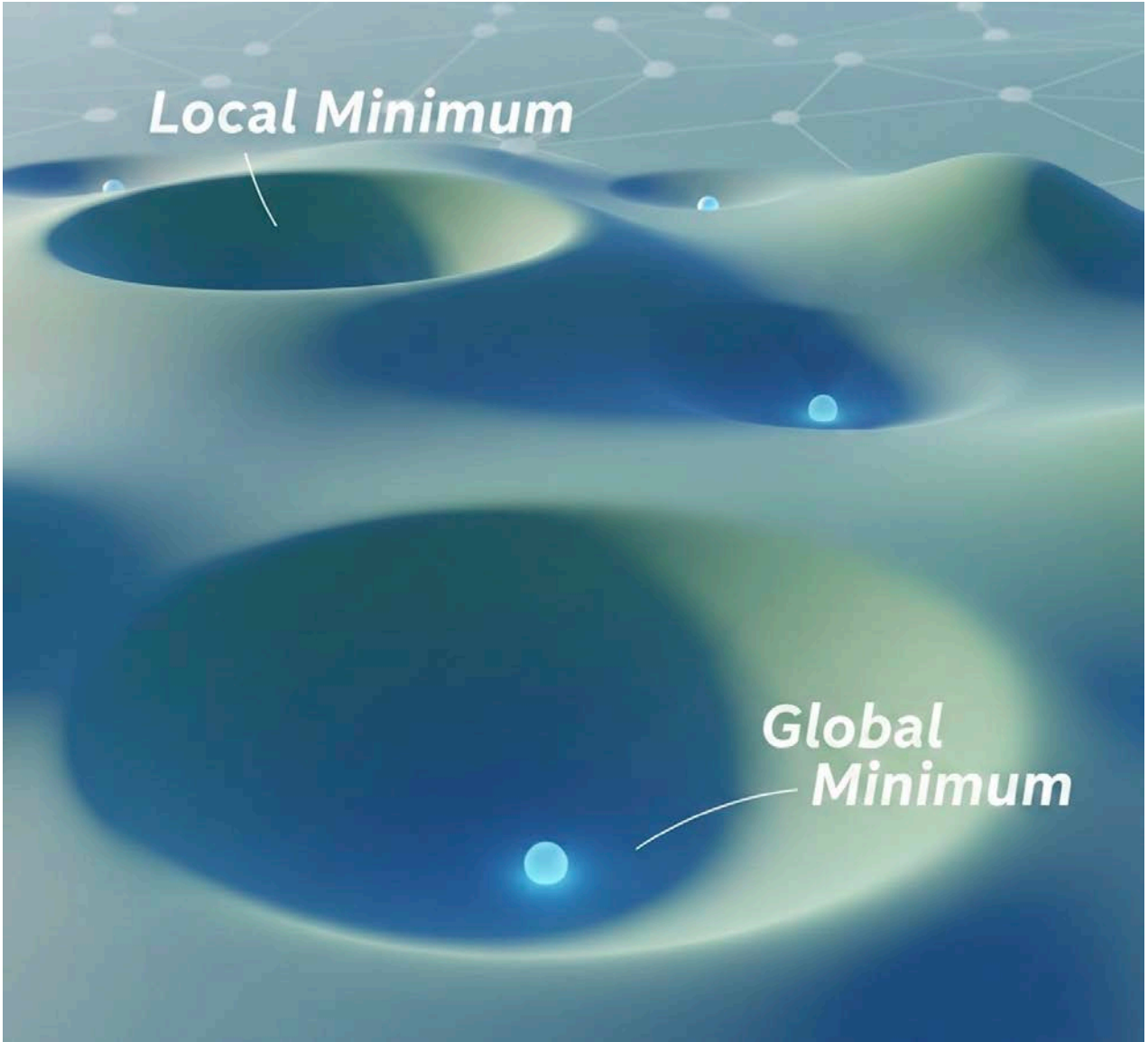
Background

Optimization problem:

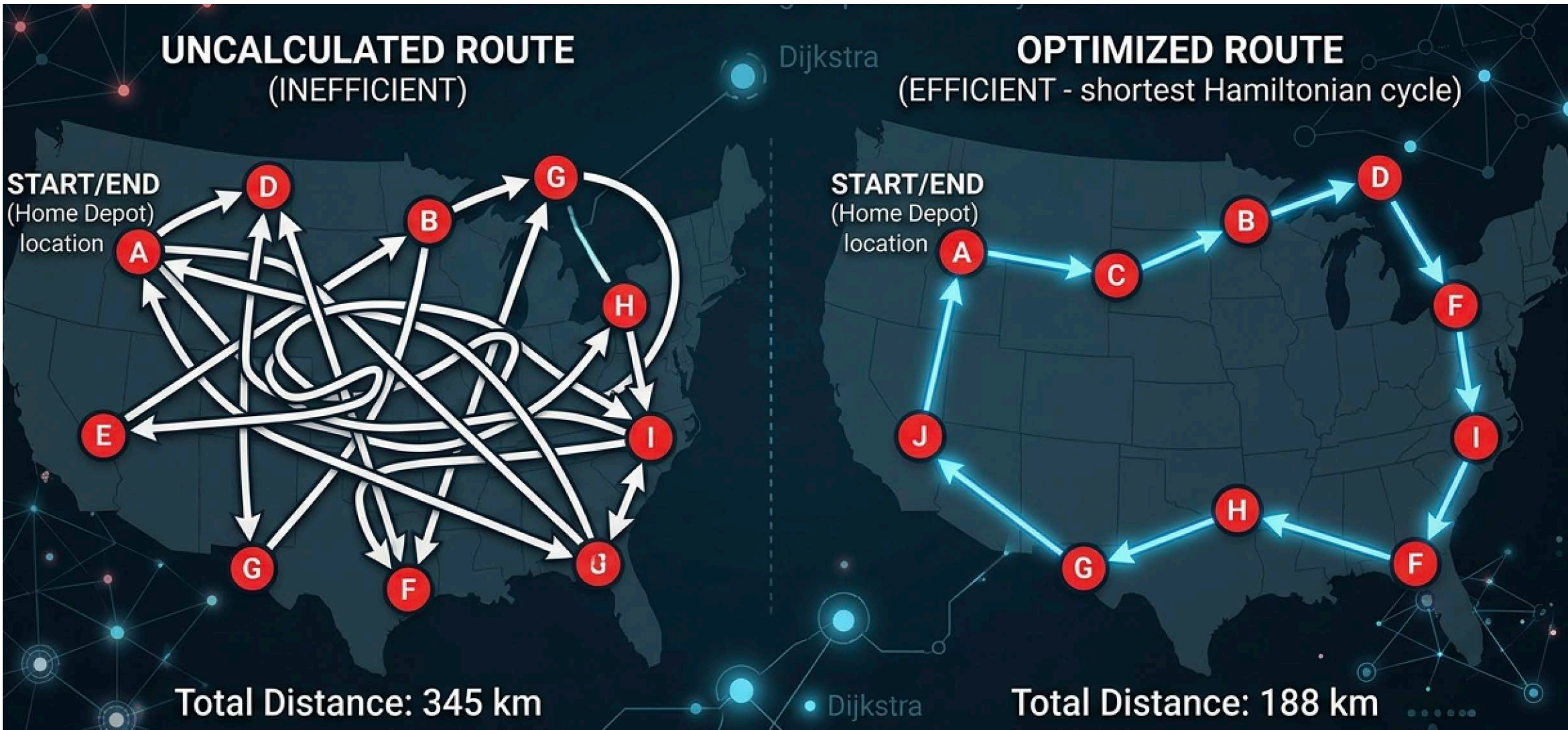
$$\max_x / \min_x f(x)$$

(Hard to solve in general)

Examples:



Finding The Ground Energy



The Traveling Salesman Problem

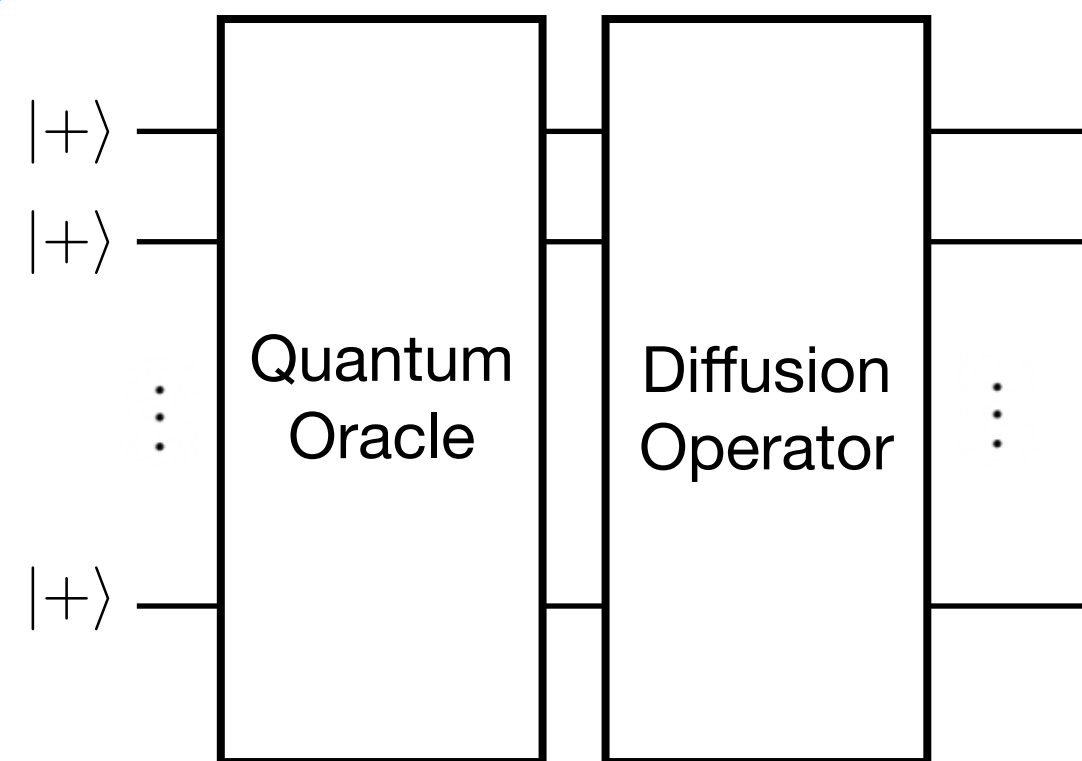
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Optimization problem:

$$\max_x / \min_x f(x)$$

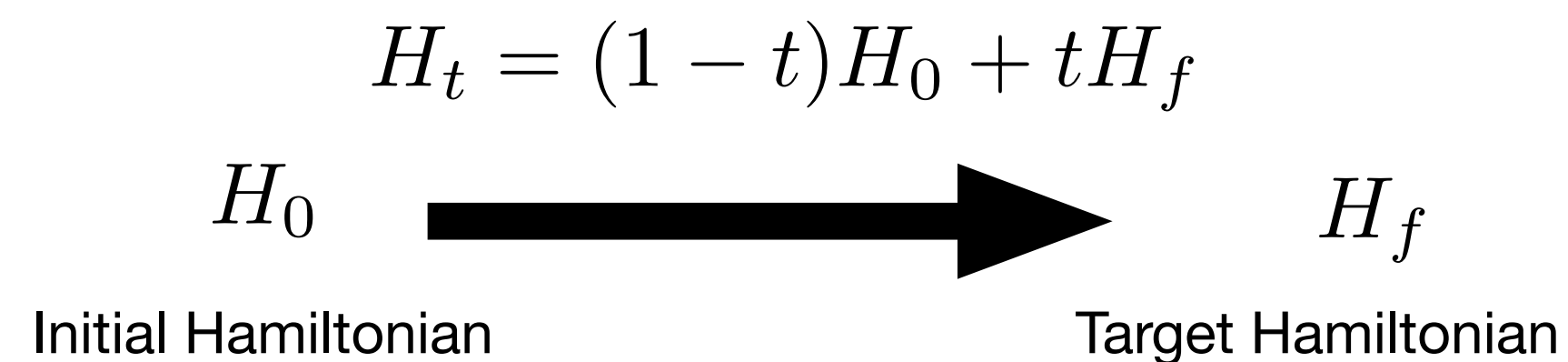
(Hard to solve in general)

Known quantum algorithms for optimization: Grover's algorithm, quantum adiabatic algorithms, VQE, QAOA ...



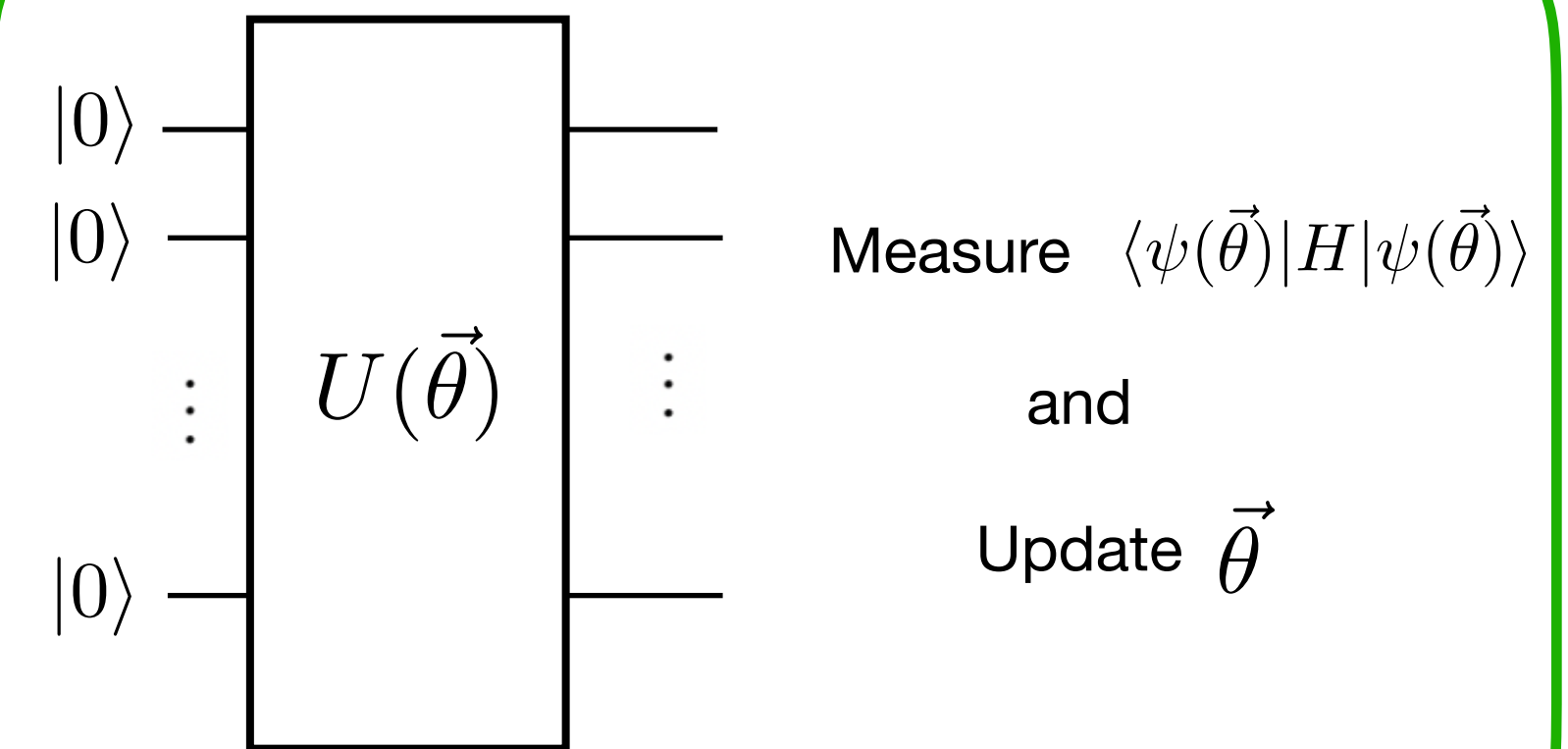
Grover's algorithm

[Grover, STOC 1996] $f(x) = 1, x = w$
 $f(x) = 0, x \neq w$



Quantum adiabatic algorithm

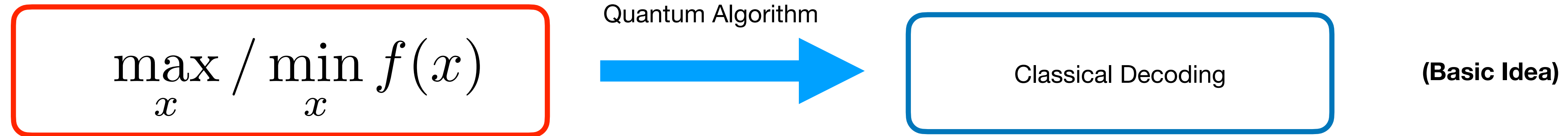
[Farhi, Goldstone, Gutmann, Sipser, *arXiv: quant-ph/0001106*]
[Aharonov, Dam, Kempe, Landau, Lloyd, Regev, FOCS 2004]



VQE



[Peruzzo, McClean, Shadbolt, Yung, Zhou, Love, Aspuru-Guzik, O'brien, Nat. Com. 2014]

Overview of DQI



Article | [Open access](#) | Published: 22 October 2025

Optimization by decoded quantum interferometry

[Stephen P. Jordan](#) , [Noah Shutty](#) , [Mary Wootters](#), [Adam Zalcman](#), [Alexander Schmidhuber](#), [Robbie King](#), [Sergei V. Isakov](#), [Tanuj Khattar](#) & [Ryan Babbush](#)

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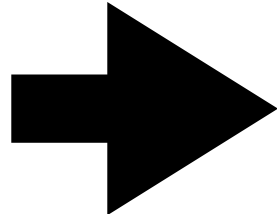
Overview of Hamiltonian DQI

Focus of today's talk

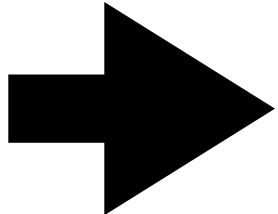
Hamiltonian

$$H = \sum_{i=1}^m c_i \text{Pauli}_i, c_i \in \mathbb{R}$$

Polynomial P



Quantum process
+
Classical decoding



$$\rho_P(H) = \frac{P(H)^2}{\text{Tr}[P(H)^2]}$$

(Basic Idea)

If the degree l is large enough, then output state $\rho_P(H) = \frac{P(H)^2}{\text{Tr}[P(H)^2]}$ is close to the Gibbs state $\frac{e^{-\beta H}}{\text{Tr}[e^{-\beta H}]}$

[Schmidhuber, Lu, Shutty, Jordan, Poremba, Quek, arXiv:2510.07913]

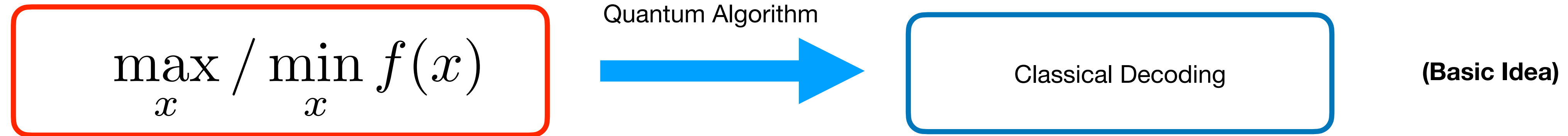
$$H = \sum_i (-1)^{v_i} \text{Pauli}_i$$

[Bu, Gu, Li, arXiv:2601.18773]

$$H = \sum_i c_i \text{Pauli}_i, c_i \in \mathbb{R}$$



Part 1: Basic Knowledge of DQI

Introduction of DQI



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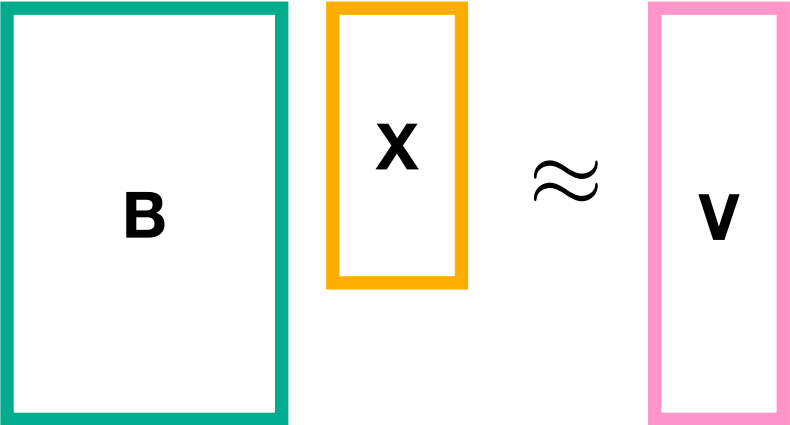
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Introduction of DQI

Example: MAX-LINSAT

Input: $B \in \mathbb{F}_2^{m \times n}$ $v \in \mathbb{F}_2^m$

Goal: Find an assignment x such that it satisfies as many constraints as possible



$$\max_{x \in \mathbb{F}_2^n} |\{i : b_i x = v_i\}| \quad b_i \text{ is the } i\text{-th row of } B$$



$$\max_{x \in \mathbb{F}_2^n} f(x), \quad f(x) = \sum_{i=1}^m (-1)^{b_i x + v_i}$$

Introduction of DQI

Example: MAX-LINSAT

Input: $B \in \mathbb{F}_2^{m \times n}$ $v \in \mathbb{F}_2^m$

Goal: $\max_{x \in \mathbb{F}_2^n} f(x), f(x) = \sum_{i=1}^m (-1)^{b_i x + v_i}$ b_i is the i -th row of B

DQI can prepare a “DQI” state

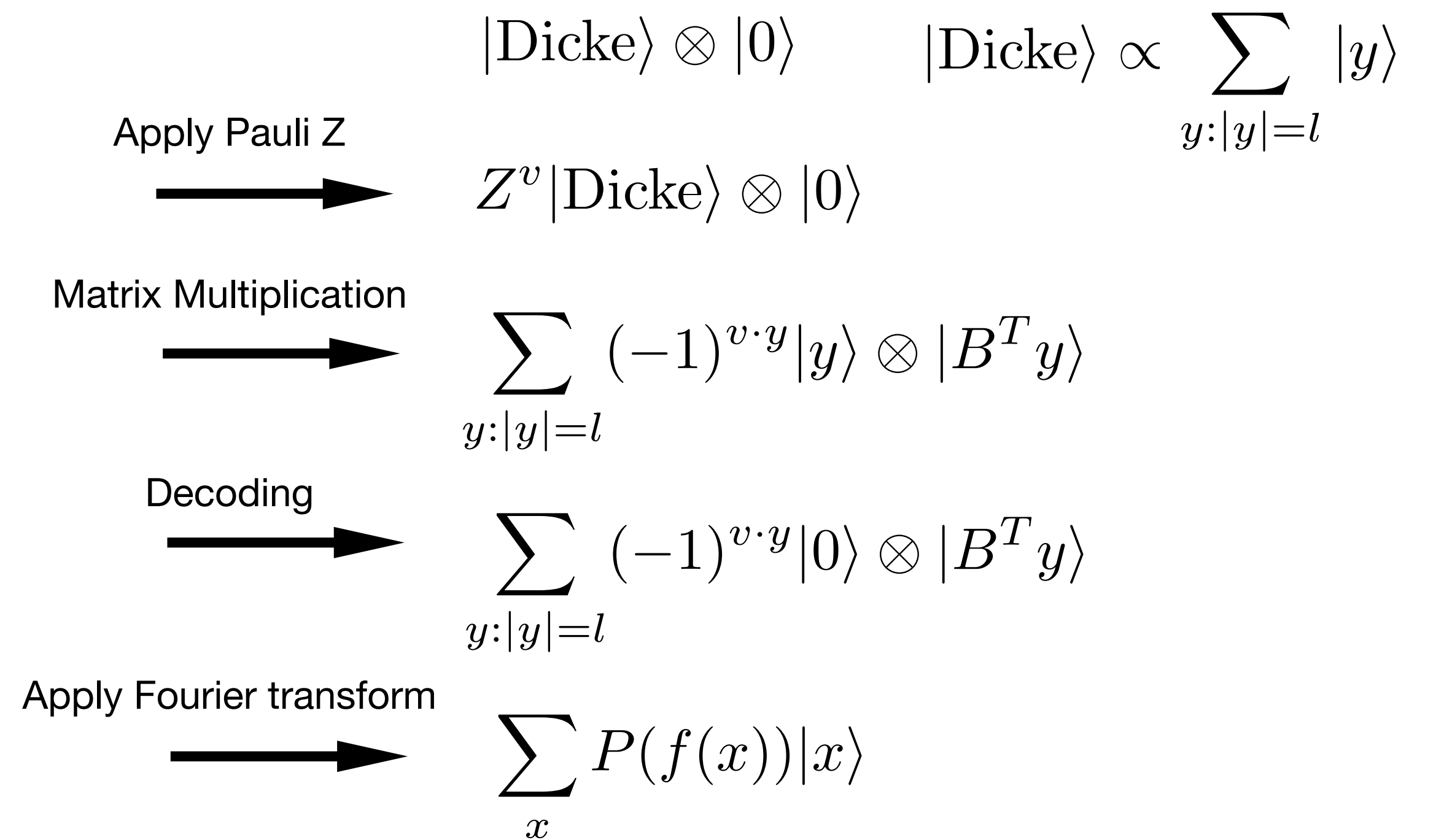
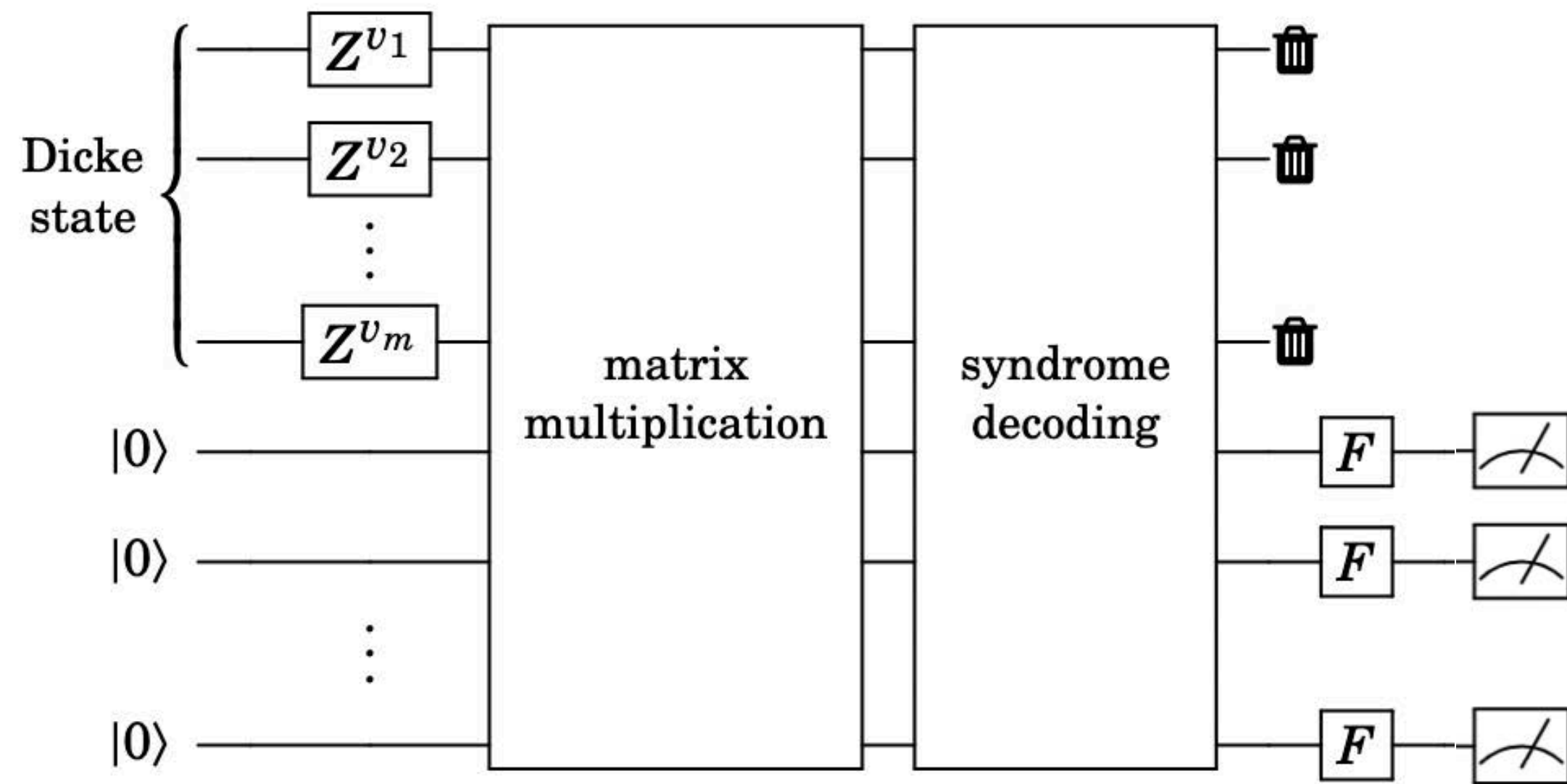
$$|P(f)\rangle \propto \sum_{x \in \mathbb{F}_2^n} P(f(x)) |x\rangle$$

P is any polynomial of degree l

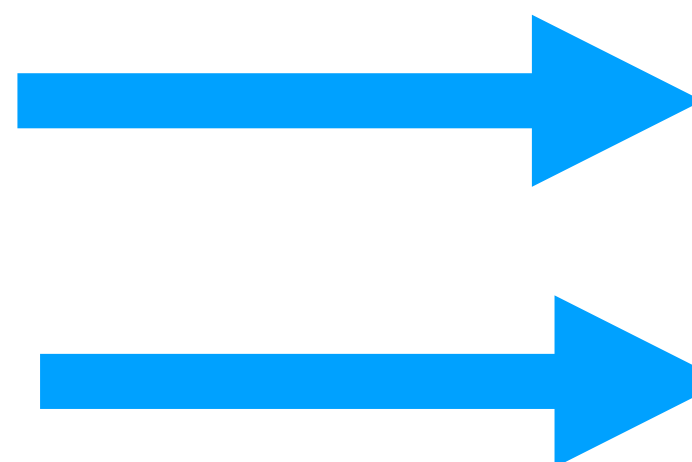
using the decoding on the code

$$C^\perp = \{y \in \mathbb{F}_2^m : B^\top y = 0\}$$

Introduction of DQI



Hence, if we recover y from $B^T y$ for $|y| \leq l$
 That is, decode 1 error for $C^\perp = \{y \in \mathbb{F}_2^m : B^T y = 0\}$

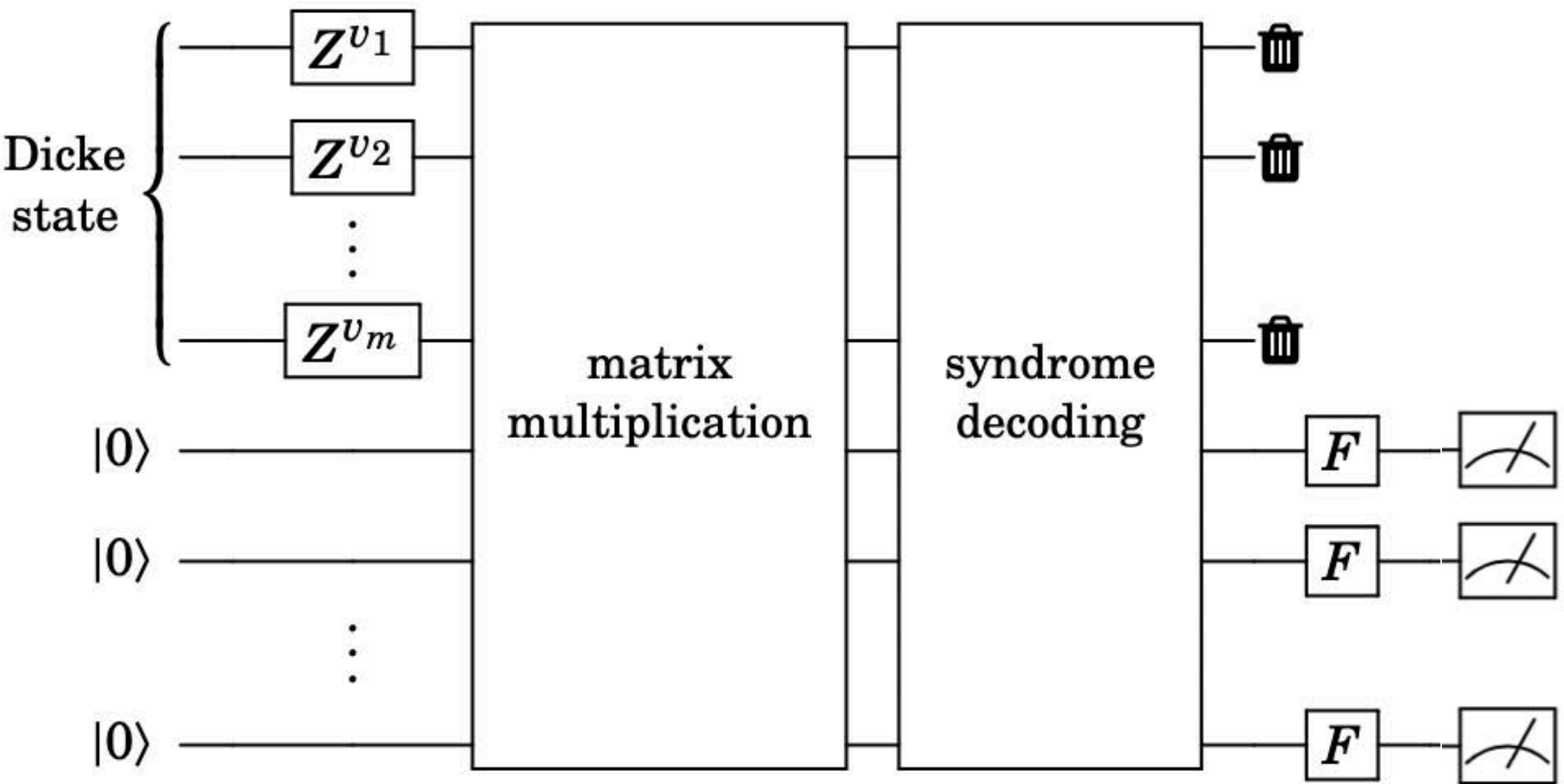


$$|P(f)\rangle \propto \sum_{x \in \mathbb{F}_2} P(f(x)) |x\rangle$$

Sampling x w.r.t probability $P(f(x))^2$

$$\langle \text{Satisfied constraints} \rangle = \sum_x f(x) P(f(x))^2$$

Introduction of DQI



Efficient decoding



Efficiency of DQI

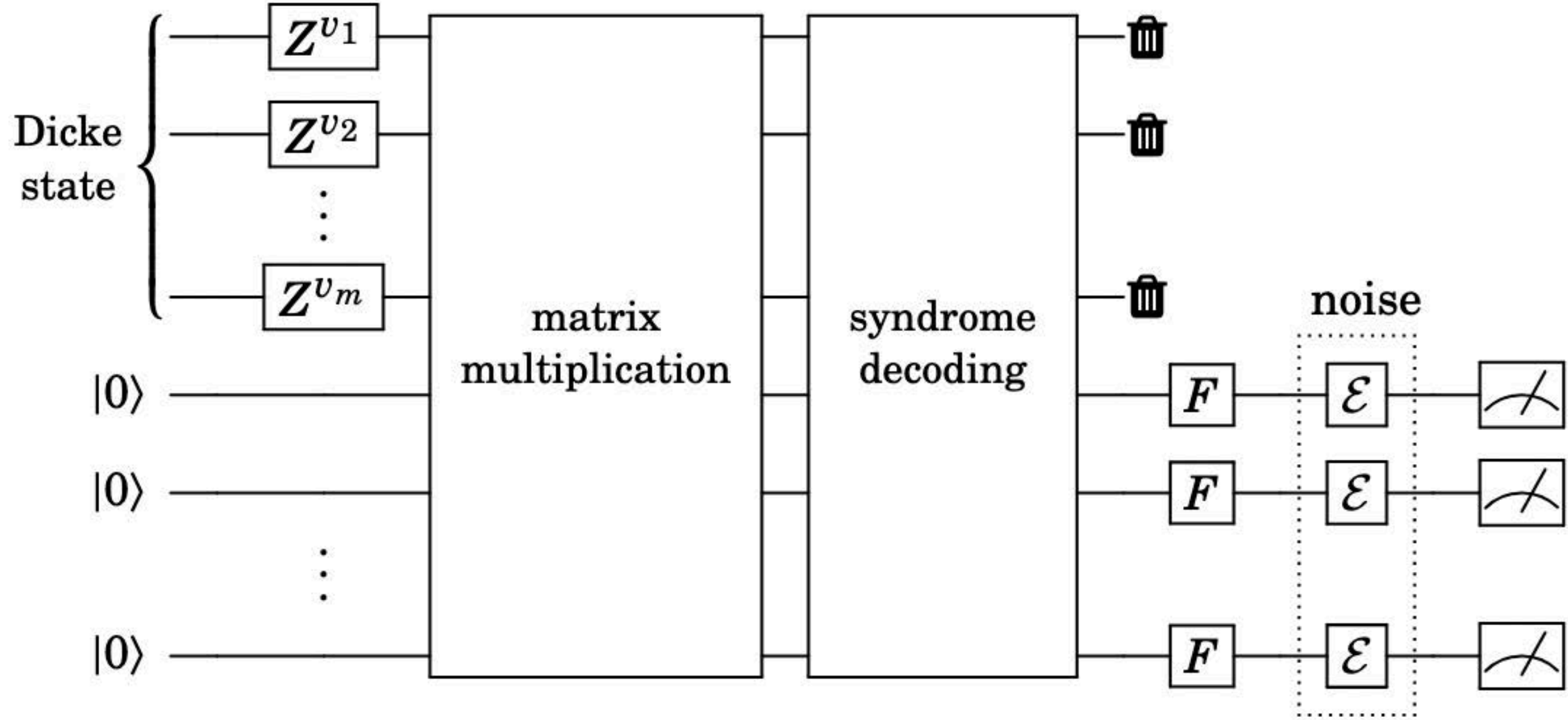
Example: Optimal Polynomial Interpolation (OPI) over \mathbb{F}_p p is a large prime number

$$B_{ij} = \gamma^{i \cdot j}, i \in \{0, \dots, p - 2\}, j \in \{0, \dots, n - 1\}, \gamma \in \mathbb{F}_p \quad \text{Reed-Solomon code}$$

There exists efficient decoding algorithm for $|y|$ less than 1/2 of the code distance (Berlekamp-Massey algorithm)

[Jordan, Shutter, Wootters, Zalcman, Schmidhuber, King, Isakov, Khattar, Babbush, Nature 2025]

Introduction of Noisy DQI



Noise: 1-qubit depolarizing noise with noise rate p

$$\mathcal{E}(X) = (1 - p)X + p \frac{\text{Tr}[X]I}{2}$$

$$\langle \text{Satisfied constraints} \rangle_{\text{Noisy}} = [\mathbb{E}_i (1 - p)^{|b_i|}] \langle \text{Satisfied constraints} \rangle_{\text{Noisless}}$$

(Exponential decay with respect to the weight of the matrix B)

[Bu, Gu, Koh, Li, Quantum Science and Technology, 2026]

From DQI to Hamiltonian DQI

DQI in Hamiltonian form

$$f(x) = \sum_{i=1}^m (-1)^{b_i x + v_i} \quad \longrightarrow \quad \text{Hamiltonian} \quad H_f = \sum_i (-1)^{v_i} Z^{b_i} \quad H_f |x\rangle = f(x) |x\rangle$$

DQI generate the sampling according to $P(f(x))^2$



Hamiltonian DQI (HDQI) [Schmidhuber, Lu, Shutter, Jordan, Poremba, Quek, *arXiv:2510.07913*]

Consider Hamiltonian $H = \sum_i (-1)^{v_i} \text{Pauli}_i \quad v_i \in \mathbb{F}_2$

HDQI aim to generate some density matrix

$$\rho_P(H) = \frac{P(H)^2}{\text{Tr}[P(H)^2]}$$

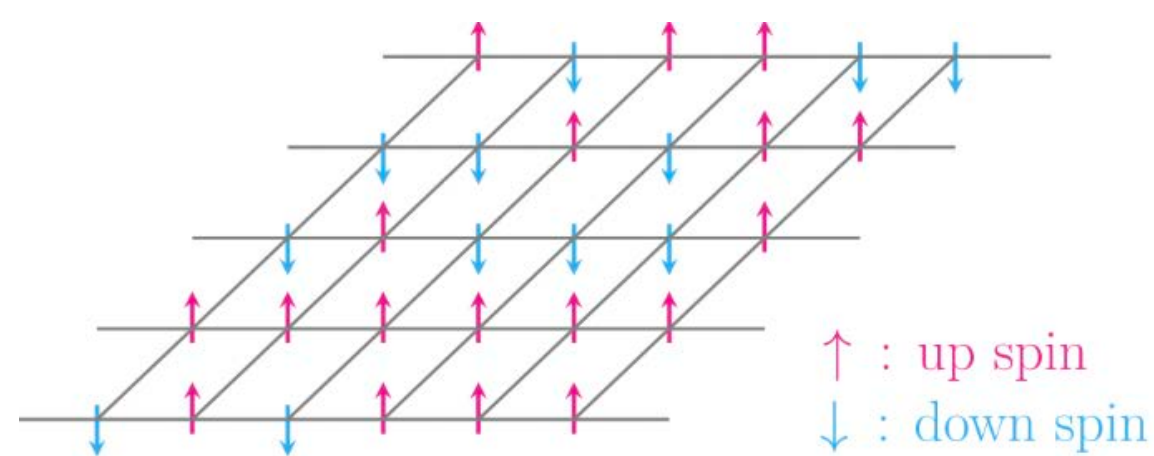
A nature question arise:
How about the physics/chemistry Hamiltonians?

If the polynomial P is a good approximation of $e^{-\beta x/2}$ \longrightarrow $\rho_P(H)$ is a good approximation of the Gibbs state

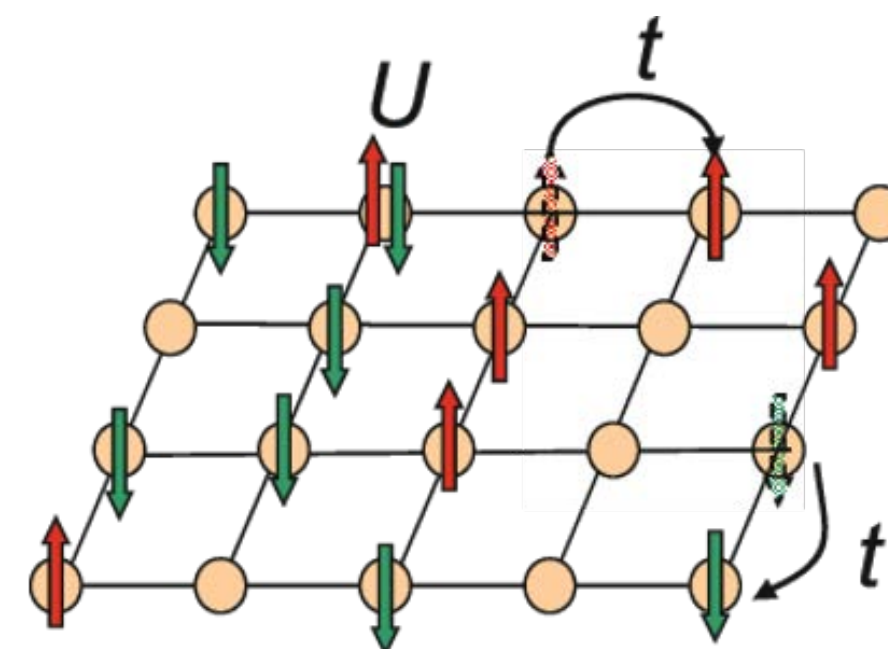
Part 2: HDQI for General Hamiltonians

Hamiltonian DQI for general Pauli Hamiltonians

Given an n-qubit Hamiltonian $H = \sum_{i=1}^m c_i P_i, c_i \in \mathbb{R}, |c_i| \leq 1$



Ising model



Hubbard model

Goal: prepare the state $\rho_P(H) = \frac{P(H)^2}{\text{Tr}[P(H)^2]}$ where P is any degree-l polynomial

Basic Knowledge of Pauli operators

Pauli operators and its symplectic representation

$$X = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}, \quad Z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}.$$

1-qubit Pauli can be written as $W(\alpha, \beta) = i^{-\alpha\beta} Z^\alpha X^\beta \quad \alpha, \beta \in \mathbb{F}_2$

$$Y = W(1, 1) = i^{-1} ZX = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}.$$

$$W(\alpha, \beta)W(\alpha', \beta') = (-1)^{\langle(\alpha, \beta), (\alpha', \beta')\rangle_s} W(\alpha', \beta')W(\alpha, \beta),$$

$$\langle(\alpha, \beta), (\alpha', \beta')\rangle_s = \alpha\beta' - \beta\alpha'$$

Basic Knowledge of Pauli operators

n-qubit Pauli can be written as $W(\vec{\alpha}, \vec{\beta}) = W(\alpha_1, \beta_1) \otimes W(\alpha_2, \beta_2) \otimes \dots \otimes W(\alpha_n, \beta_n)$,

$$W(\vec{\alpha}, \vec{\beta})W(\vec{\alpha}', \vec{\beta}') = (-1)^{\langle(\vec{\alpha}, \vec{\beta}), (\vec{\alpha}', \vec{\beta}')\rangle_s} W(\vec{\alpha}', \vec{\beta}')W(\vec{\alpha}, \vec{\beta}),$$

So the vector $(\vec{\alpha}, \vec{\beta}) \in \mathbb{F}_2^{2n}$ is called its symplectic representation for a Pauli operator P , denoted by $\text{symp}(P)$

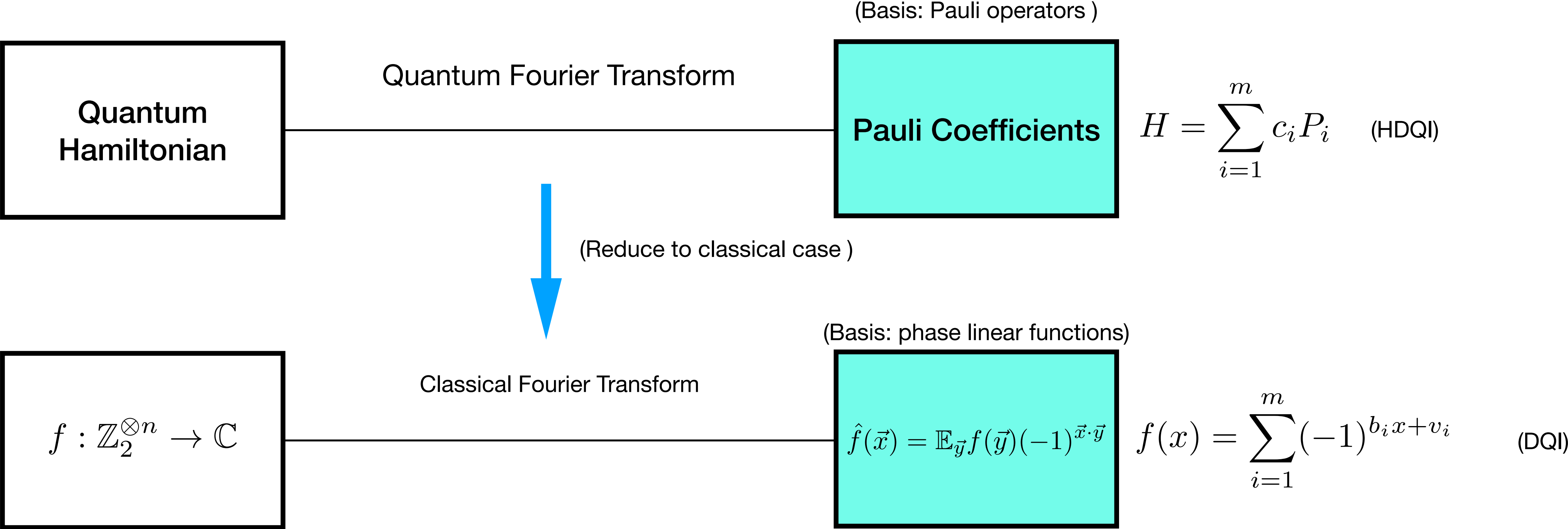
Moreover, $W(\vec{\alpha}, \vec{\beta})W(\vec{\alpha}', \vec{\beta}') = i^{\langle(\vec{\alpha}, \vec{\beta}), (\vec{\alpha}', \vec{\beta}')\rangle_s} W(\vec{\alpha} + \vec{\alpha}', \vec{\beta} + \vec{\beta}')$

➔ Group homomorphism: $\text{symp}(PQ) = \text{symp}(P) \oplus \text{symp}(Q)$

$$|\text{symp}(P)\rangle \begin{array}{c} \leftarrow \text{CNOT, H} \rightarrow \\ \text{CNOT, H} \end{array} P \otimes I |\text{Bell}_n\rangle$$

Pauli operators (via Fourier analysis)

(Pauli operators: ONB for $B(\mathcal{H}^{\otimes n})$ with respect to $\langle A, B \rangle = \frac{1}{2^n} \text{Tr}[A^\dagger B]$)



Remark: Pauli operators are taken as the “quantum polynomials” of degree 1 in quantum higher-order Fourier analysis [Bu, Gu, Jaffe, PNAS 2025]

Hamiltonian DQI for general Pauli Hamiltonians

Using symplectic representation

$$H = \sum_{i=1}^m c_i P_i \quad \longrightarrow \quad B^\top = \begin{bmatrix} | & | & & | \\ \text{symp}(P_1) & \text{symp}(P_2) & \cdots & \text{symp}(P_m) \\ | & | & & | \end{bmatrix} \in \mathbb{F}_2^{2n \times m}$$

(Classical Linear Code) $C^\perp = \{y \in \mathbb{F}_2^m : B^\top y = 0\}$

Assume there is an efficient and perfect decoder

$$D_H^{(l)} |y\rangle |B^\top y\rangle = |0\rangle |B^\top y\rangle, \forall |y| \leq l$$

For example, if the columns are linear independent, then using Gaussian elimination to decode

Commuting case

Given: $H = \sum_{i=1}^m c_i P_i$ Polynomial P

Goal: $\rho_P(H) = \frac{P(H)^2}{\text{Tr}[P(H)^2]}$

Step 0: rewrite the polynomial P(H)

For any degree- l polynomial $P(x) = \sum_{j=0}^l a_j x^j$ consider $P(H) = \sum_{j=0}^l a_j H^j$

It can be rewritten as $P(H) = \sum_{\mathbf{y} \in \mathbb{F}_2^m} w_{\mathbf{y}} P^{\mathbf{y}}$ $P^{\mathbf{y}} = P_1^{y_1} \cdot P_2^{y_2} \cdot \dots \cdot P_m^{y_m}$

$$w_{\mathbf{y}} = \sum_{j=0}^l a_j j! \sum_{\mu \in \mathbb{Z}_+^m : |\mu|=j, \mu \equiv \mathbf{y} \pmod{2}} \frac{c^{\mu}}{\mu!} \quad c^{\mu} = c_1^{\mu_1} \cdot c_2^{\mu_2} \cdot \dots \cdot c_m^{\mu_m} \quad \mu! = \mu_1! \cdot \mu_2! \cdot \dots \cdot \mu_m!$$

Commuting case

Step 1: Prepare a reference state to encode $P(H)$ $|R^l(H)\rangle = \frac{1}{\mathcal{N}} \sum_{\mathbf{y} \in \mathbb{F}_2^m} w_{\mathbf{y}} |\mathbf{y}\rangle$ Recall that $P(H) = \sum_{\mathbf{y} \in \mathbb{F}_2^m} w_{\mathbf{y}} P^{\mathbf{y}}$

Step 2: Apply control-Pauli operators between the reference state and Bell states on $2n$ qubits

$$\frac{1}{\mathcal{N}} \sum_{\mathbf{y} \in \mathbb{F}_2^m} w_{\mathbf{y}} |\mathbf{y}\rangle \otimes (P^{\mathbf{y}} \otimes I) |\text{Bell}_n\rangle \longleftrightarrow \frac{1}{\mathcal{N}} \sum_{\mathbf{y} \in \mathbb{F}_2^m} w_{\mathbf{y}} |\mathbf{y}\rangle \otimes |\text{Symp}(P^{\mathbf{y}})\rangle \quad \text{Symp}(P^{\mathbf{y}}) = B^T \mathbf{y}$$

$$B^T = \begin{bmatrix} | & | & \cdots & | \\ \text{symp}(P_1) & \text{symp}(P_2) & \cdots & \text{symp}(P_m) \\ | & | & \cdots & | \end{bmatrix}$$

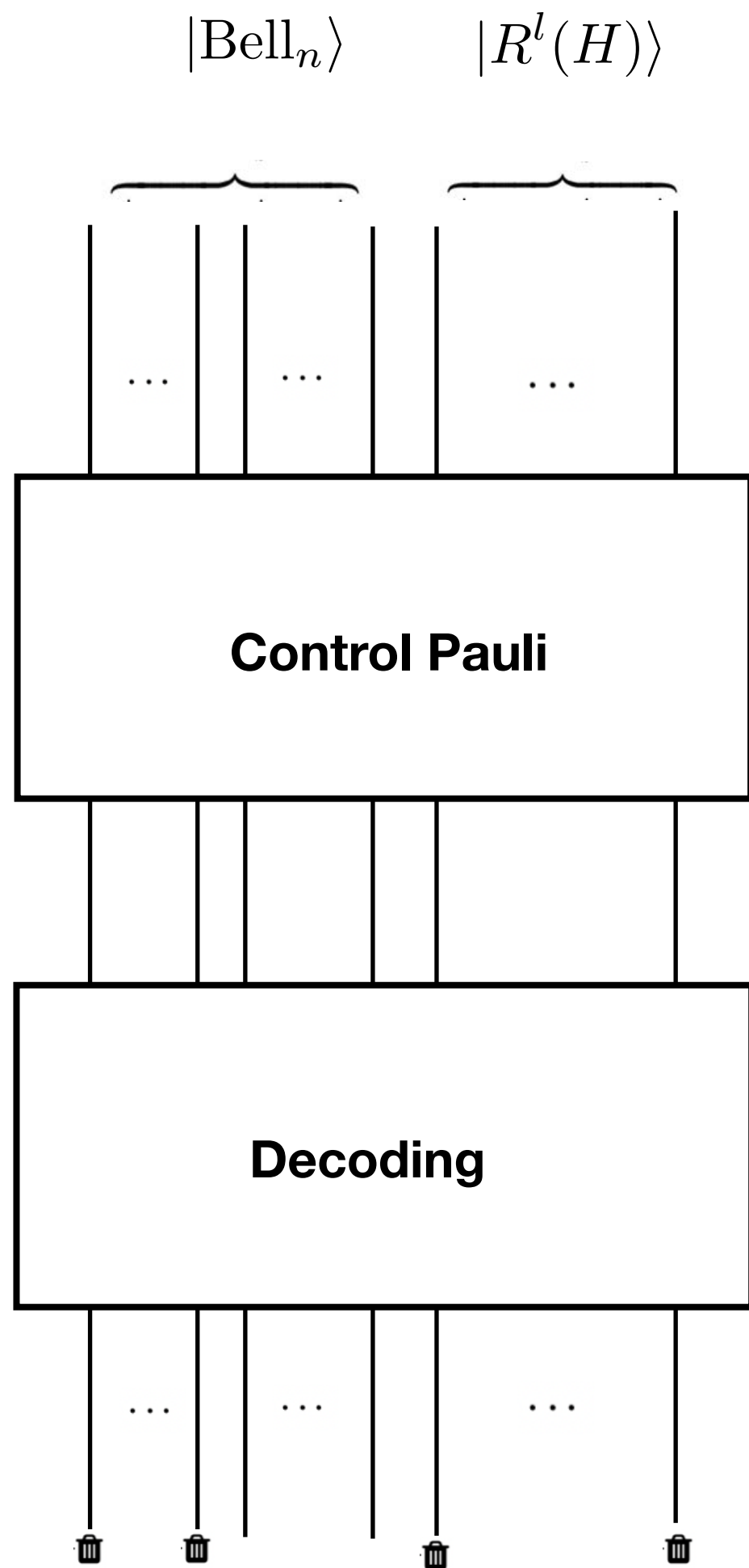
Step 3: Decoding

$$\frac{1}{\mathcal{N}} \sum_{\mathbf{y} \in \mathbb{F}_2^m} w_{\mathbf{y}} |\mathbf{0}\rangle \otimes |\text{Symp}(P^{\mathbf{y}})\rangle \longleftrightarrow \frac{1}{\mathcal{N}} \sum_{\mathbf{y} \in \mathbb{F}_2^m} w_{\mathbf{y}} (P^{\mathbf{y}} \otimes I) |\text{Bell}_n\rangle = P(H) \otimes I |\text{Bell}_n\rangle$$

Step 4: Take partial trace

$$P(H) \otimes I |\text{Bell}_n\rangle \longrightarrow \rho_P(H) = \frac{P(H)^2}{\text{Tr}[P(H)^2]}$$

Commuting case



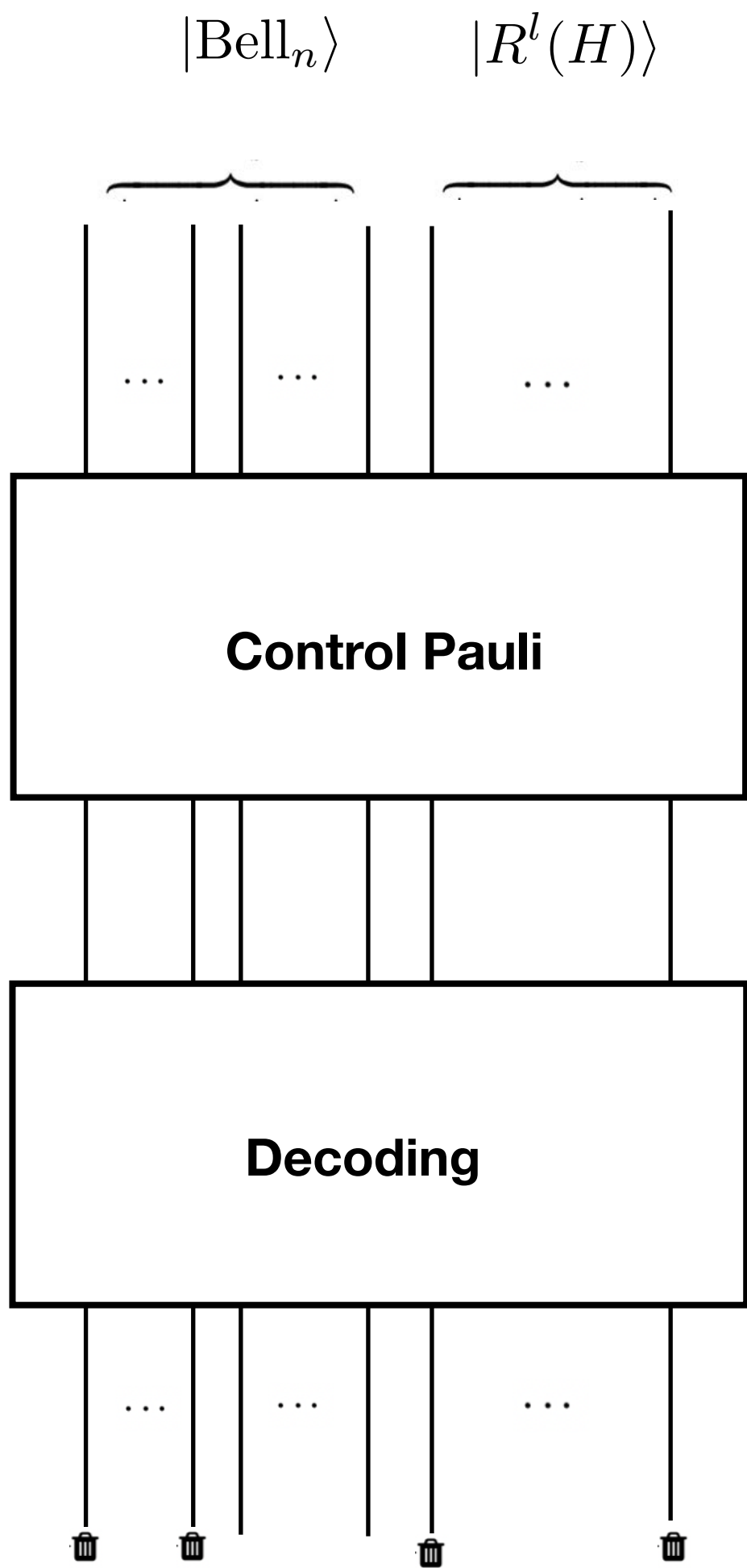
$$\begin{aligned}
 & |R^l(H)\rangle \otimes |\text{Bell}_n\rangle \\
 & \downarrow \\
 & \frac{1}{\mathcal{N}} \sum_{\mathbf{y} \in \mathbb{F}_2^m} w_{\mathbf{y}} |\mathbf{y}\rangle \otimes (P^{\mathbf{y}} \otimes I) |\text{Bell}_n\rangle \\
 & = \frac{1}{\mathcal{N}} \sum_{\mathbf{y} \in \mathbb{F}_2^m} w_{\mathbf{y}} |\mathbf{y}\rangle \otimes |\text{Symp}(P^{\mathbf{y}})\rangle \\
 & \downarrow \\
 & \frac{1}{\mathcal{N}} \sum_{\mathbf{y} \in \mathbb{F}_2^m} w_{\mathbf{y}} |0\rangle \otimes |\text{Symp}(P^{\mathbf{y}})\rangle = P(H) \otimes I |\text{Bell}_n\rangle \\
 & \downarrow \\
 & \rho_P(H) = \frac{P(H)^2}{\text{Tr}[P(H)^2]}
 \end{aligned}$$

$$P(H) = \sum_{\mathbf{y} \in \mathbb{F}_2^m} w_{\mathbf{y}} P^{\mathbf{y}}$$

$$\text{Symp}(P^{\mathbf{y}}) = B^T \mathbf{y}$$

$$B^T = \begin{bmatrix} | & | & & | \\ \text{symp}(P_1) & \text{symp}(P_2) & \cdots & \text{symp}(P_m) \\ | & | & & | \end{bmatrix}$$

Commuting case



$$|R^l(H)\rangle \otimes |\text{Bell}_n\rangle$$

$$\downarrow$$

$$\frac{1}{\mathcal{N}} \sum_{\mathbf{y} \in \mathbb{F}_2^m} w_{\mathbf{y}} |\mathbf{y}\rangle \otimes (P^{\mathbf{y}} \otimes I) |\text{Bell}_n\rangle$$

$$= \frac{1}{\mathcal{N}} \sum_{\mathbf{y} \in \mathbb{F}_2^m} w_{\mathbf{y}} |\mathbf{y}\rangle \otimes |\text{Symp}(P^{\mathbf{y}})\rangle$$

$$\downarrow$$

$$\frac{1}{\mathcal{N}} \sum_{\mathbf{y} \in \mathbb{F}_2^m} w_{\mathbf{y}} |0\rangle \otimes |\text{Symp}(P^{\mathbf{y}})\rangle = P(H) \otimes I |\text{Bell}_n\rangle$$

$$\downarrow$$

$$\rho_P(H) = \frac{P(H)^2}{\text{Tr}[P(H)^2]}$$

$$P(H) = \sum_{\mathbf{y} \in \mathbb{F}_2^m} w_{\mathbf{y}} P^{\mathbf{y}}$$

$$\text{Symp}(P^{\mathbf{y}}) = B^T \mathbf{y}$$

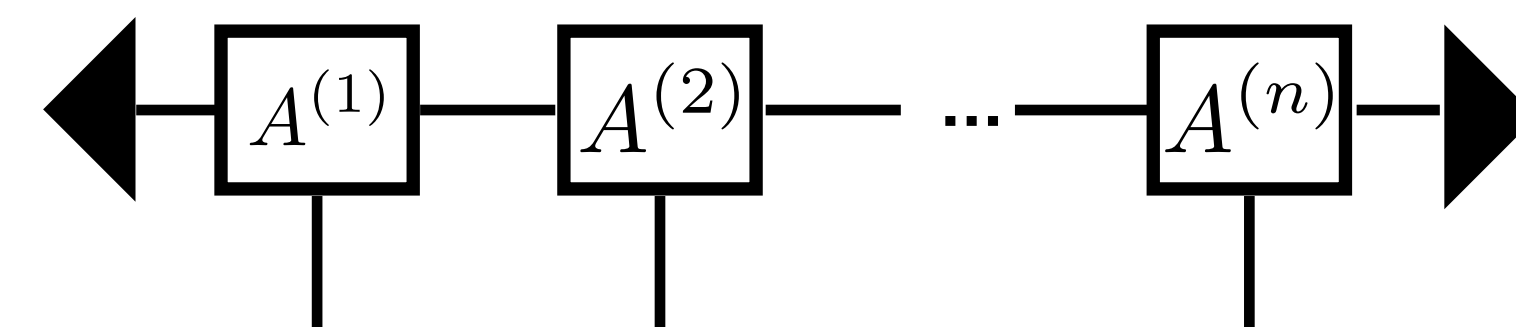
$$B^T = \begin{bmatrix} | & | & & | \\ \text{symp}(P_1) & \text{symp}(P_2) & \dots & \text{symp}(P_m) \\ | & | & & | \end{bmatrix}$$

A nature question arise:
can the reference state be efficiently prepared

Efficient preparation of reference state

Result: The reference state is an Matrix Product State (MPS) with bond dimension $l+1$

$$|R^l(H)\rangle = \sum_{y \in \mathbb{F}_2^m} v_L^\top A^{(1)}(y_1) \cdots A^{(m)}(y_m) v_R |y\rangle$$

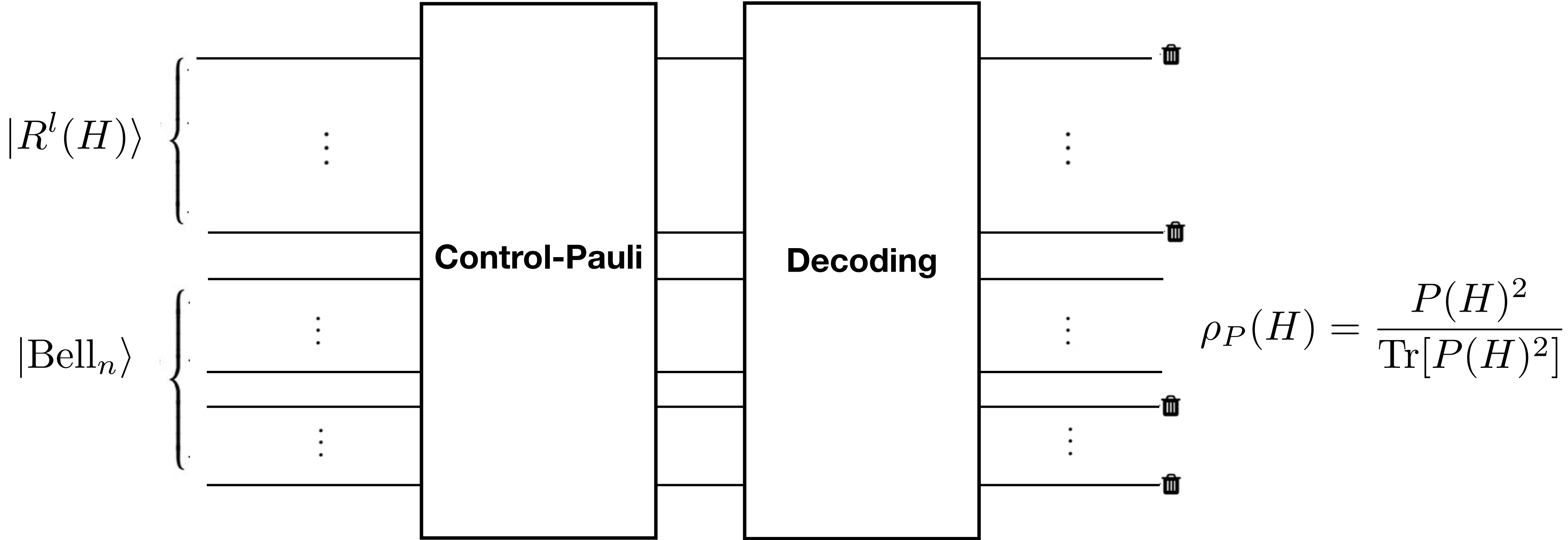


where $v_L = \left(\frac{1}{\mathcal{N}}, 0, \dots, 0\right)^\top$, $v_R = (a_0 \cdot 0!, a_1 \cdot 1!, \dots, a_l \cdot l!)^\top$

$$A_0^{(k)} = \begin{pmatrix} 1 & 0 & \frac{c_k^2}{2!} & 0 & \frac{c_k^4}{4!} & \dots & \dots \\ 0 & 1 & 0 & \frac{c_k^2}{2!} & 0 & \ddots & \vdots \\ 0 & 0 & 1 & 0 & \frac{c_k^2}{2!} & \ddots & \frac{c_k^4}{4!} \\ \vdots & \vdots & \ddots & \ddots & \ddots & \ddots & 0 \\ 0 & 0 & \dots & 0 & 1 & 0 & \frac{c_k^2}{2!} \\ 0 & 0 & \dots & 0 & 0 & 1 & 0 \\ 0 & 0 & \dots & 0 & 0 & 0 & 1 \end{pmatrix}_{(l+1) \times (l+1)}, \quad A_1^{(k)} = \begin{pmatrix} 0 & \frac{c_k}{1!} & 0 & \frac{c_k^3}{3!} & 0 & \dots & \dots \\ 0 & 0 & \frac{c_k}{1!} & 0 & \frac{c_k^3}{3!} & \ddots & \vdots \\ 0 & 0 & 0 & \frac{c_k}{1!} & 0 & \ddots & 0 \\ \vdots & \vdots & \ddots & \ddots & \ddots & \ddots & \frac{c_k^3}{3!} \\ 0 & 0 & \dots & 0 & 0 & 0 & 0 \\ 0 & 0 & \dots & 0 & 0 & 0 & \frac{c_k}{1!} \\ 0 & 0 & \dots & 0 & 0 & 0 & 0 \end{pmatrix}_{(l+1) \times (l+1)}$$

 The reference state (MPS) can be prepared in time $O(m \cdot \text{poly}(l))$

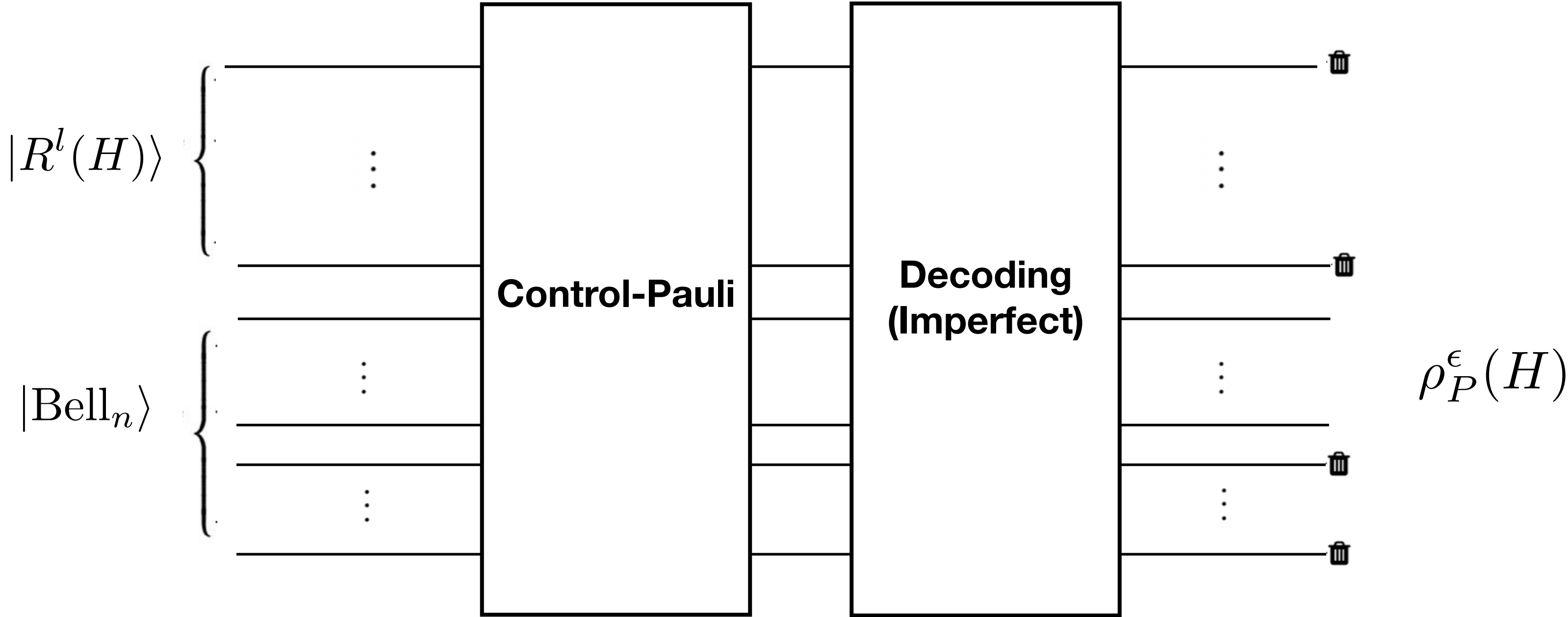
Commuting case



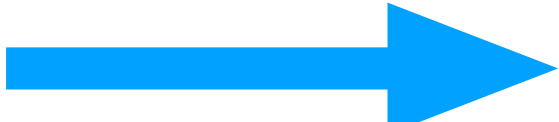
$|R^l(H)\rangle$ can be prepared efficiently + Efficient decoding

Efficiently prepare $\rho_P(H)$

Commuting case



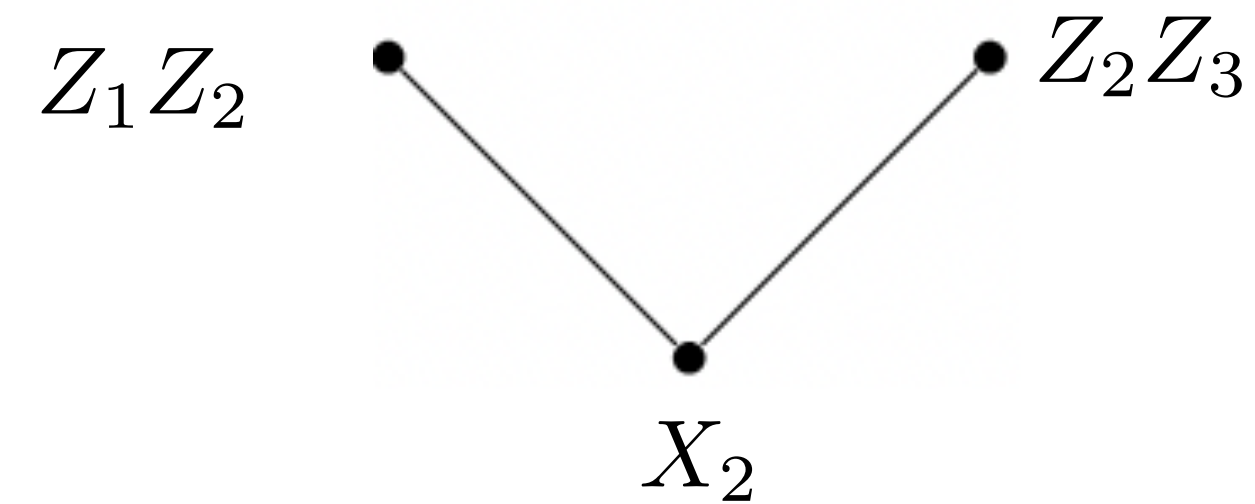
Imperfect Decoding: $D_H^{(l,\epsilon)} |y\rangle |B^\top y\rangle = \sum_{y'} \sqrt{p(y'|B^\top y)} |y \oplus y'\rangle |B^\top y\rangle, \forall |y| \leq l \quad p(y|B^\top y) \geq 1 - \epsilon$

 $\|\rho_P^\epsilon(H) - \rho_P(H)\|_1 \leq 2\sqrt{\epsilon}$ (Robustness to imperfect decoding)

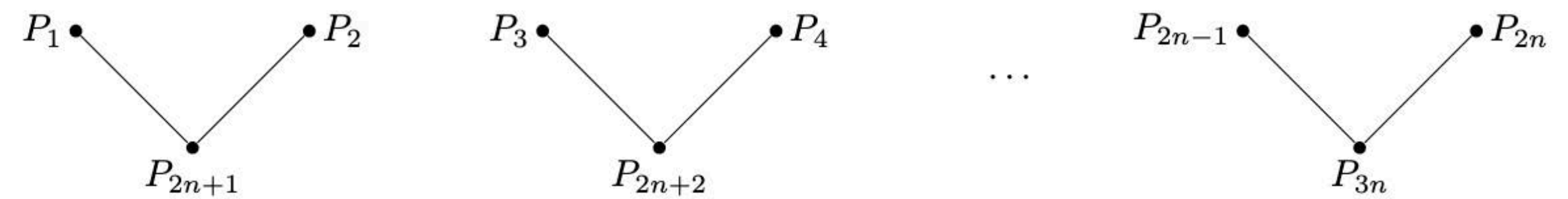
Noncommuting case

We need introduce the commutation graph based on the commute and anti-commute relations of the Pauli operators

E.g. $H = Z_1 Z_2 + Z_2 Z_3 + g X_2$

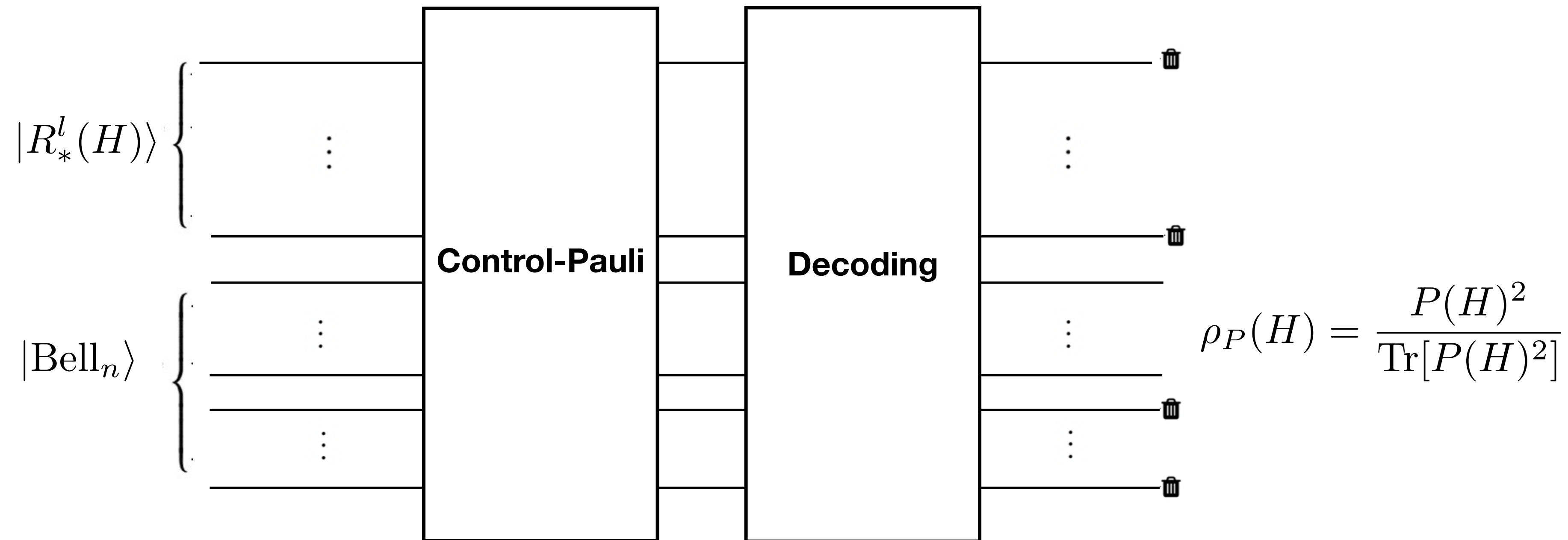


$$H = \sum_{i=1}^{2n} Z_i Z_{i+1} + g \sum_{i=1}^n X_{2i},$$



$$P_i = \begin{cases} Z_i Z_{i+1} & \text{for } 1 \leq i \leq 2n, \\ X_{2(i-2n)} & \text{for } 2n+1 \leq i \leq 3n, \end{cases}$$

Noncommuting case



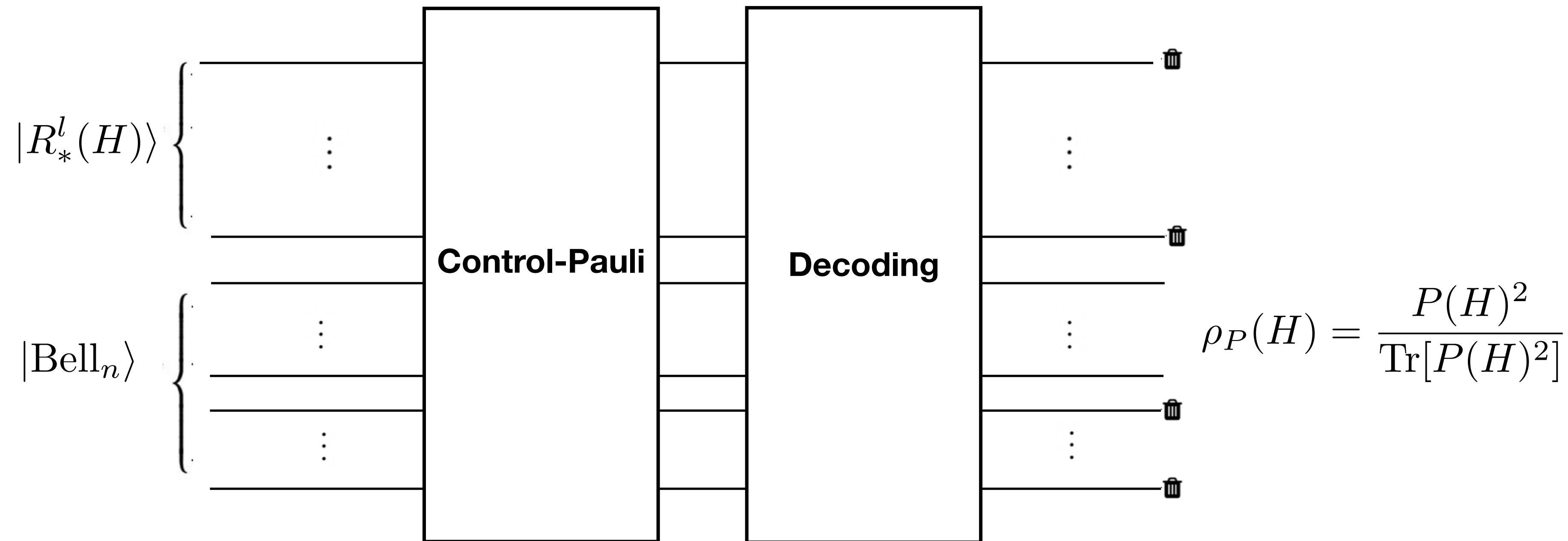
In the noncommuting case, the reference state is quite complicated

However, we can still express it into an MPS, and time complexity for preparation is

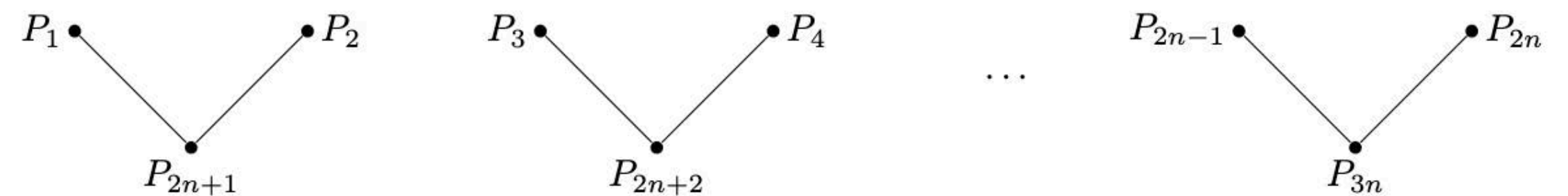
$$O(m \cdot \text{poly}(l) \cdot \exp(\mathcal{M}))$$

\mathcal{M} : maximal size of the connected components in the commutation graph of H

Noncommuting case



For example,
$$H = \sum_{i=1}^{2n} Z_i Z_{i+1} + g \sum_{i=1}^n X_{2i},$$

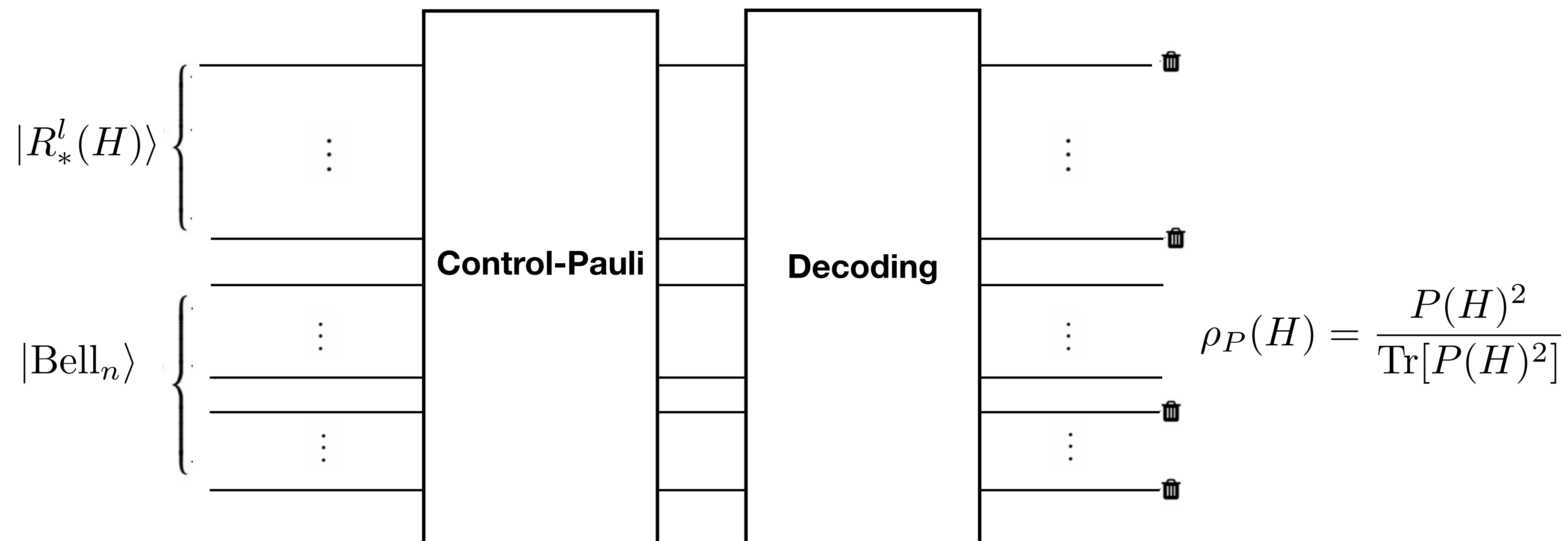


In this example, the maximal size of the connected components is 3, i.e., $\mathcal{M} = 3$

So the time complexity to prepare $|R_*^l(H)\rangle$ is $O(m \cdot \text{poly}(l))$

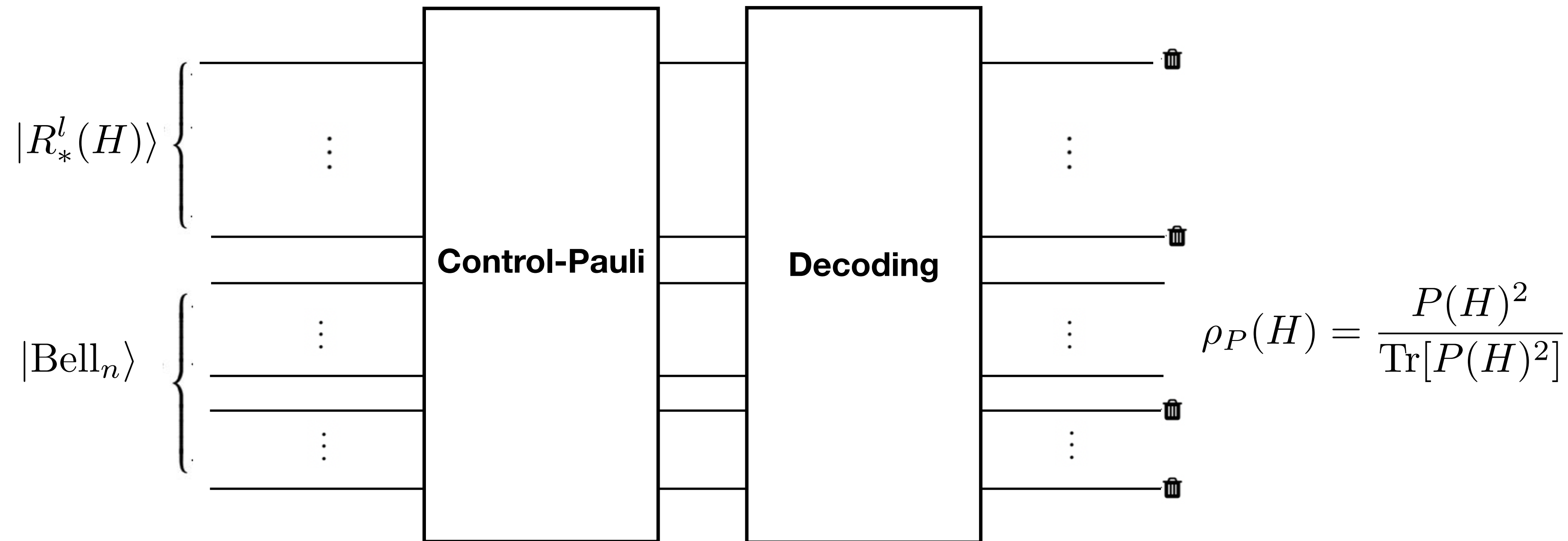
Details in [Bu, Gu, Li, arXiv:2601.18773]

Application: approximating Gibbs state



If the degree l is large enough, then output state $\rho_P(H) = \frac{P(H)^2}{\text{Tr}[P(H)^2]}$ is close to the Gibbs state $\frac{e^{-\beta H}}{\text{Tr}[e^{-\beta H}]}$

Application: approximating Gibbs state



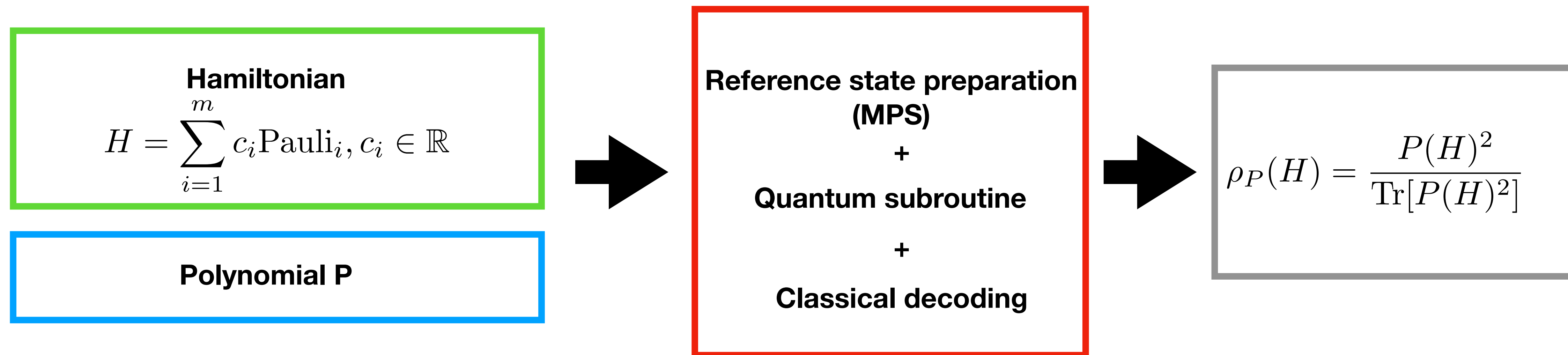
For example, $H = \sum_{i=1}^{2n} Z_i Z_{i+1} + g \sum_{i=1}^n X_{2i}$, we can choose some polynomial P with degree $l \leq 1.12 (2 + |g|)\beta n + 0.648 \ln \frac{2}{\delta}$

$$\|\rho_P(H) - e^{-\beta H} / \text{Tr}[e^{-\beta H}]\|_1 \leq \delta$$

And the total running time is $\text{poly}(l, n) = \text{poly}((2 + |g|)\beta, \ln \frac{1}{\delta}, n)$

Summary and outlook

In this work, we introduce a quantum algorithm to generate a function calculus on Hamiltonian $H = \sum_{i=1}^m c_i \text{Pauli}_i, c_i \in \mathbb{R}, |c_i| \leq 1$



Future directions: (1) More efficient algorithm for noncommuting case

(2) Other applications

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Thank You!

