

My Adventures in Quantumland

Man-Duen Choi

choi@math.toronto.edu

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***MATHEMATICAL PICTURE
LANGUAGE SEMINAR***

Outline

➤ ***What is a Quantum Computer?***

- Qubits

- density matrices

➤ ***The **Taming** of Shrews***

- tensor products

- Quantum Entanglements

- Completely positive linear maps

➤ ***Old and **New** Results***

Units of Information -- Bits vs Qubits

- A **bit** (binary integer) is the base of conventional computer memory.
- **1-bit** is read as either a zero or a one with **probability** in the real **interval** $[0, 1]$.
- A **3-bit** corresponds to an element in $\{0, 1\} \times \{0, 1\} \times \{0, 1\} = 2^3$, as vertices of a cube;

but very WRONG to have a **cube** for **probability**!
- When $n = 40$, we get $2^{40} = \text{tera}$

- To get a setting of a possible non-commutative generalization, we associate each **1-bit** with a rank-1 diagonal 2×2 projection matrix, i.e.,

$$\begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} \quad \text{or} \quad \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix}$$

- Each **3-bit** corresponds to a rank-1 diagonal 8×8 projection matrix, as the tensor product of three 2×2 matrices where each is

$$\begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} \quad \text{or} \quad \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix}$$

Thus there are eight **3-bits** located in an 8-dimensional space.

- Apparently, nobody in computer science mentioned of diagonal matrices and tensor products.

- A **qubit** (quantum bit) is a unit of quantum computer memory.
- Mathematically, each **1-qubit** is regarded as an element in

$$S^2 \simeq \left\{ \frac{1}{2} \begin{pmatrix} 1-x & y+iz \\ y-iz & 1+x \end{pmatrix} \quad \text{with } x^2 + y^2 + z^2 = 1 \right\}$$

- = {all 2 X 2 rank-1 projection matrices}
- = {all vector states acting on \mathbf{C}^2 }
- = {one-dimensional complex linear subspaces of \mathbf{C}^2 }
- = {*special* two-dimensional real subspaces of \mathbf{R}^4 }

- Physically, a **1-qubit** is a superposition of the spherical surface (called the **Bloch sphere**), because an “electron” can move freely to any direction from the origin of \mathbf{R}^3 .
- Thus, S^2 need not be a material surface.
- There are **uncountably many** 1-qubits, to make S^2 **symmetry** with **continuity** for **approximation**, while 1-bits are just the north pole and the south pole.

What on earth does S^2 mean?

- Think of geography and physics and philosophy, instead of set theory and computers.
- It means of measure theory (as length / area / volume) and continuity and dimension and analog. Always uncountably infinite points (beyond the capacity of any conventional computer memory).
- It goes along with human memory, which could be transcendental and sensible and sensational and sentimental.

- However, S^2 is a mathematical simple object. The combinatorial effect of S^2 symmetry is not comparable to 2^n when $n > 30$. So, a 1-qubit computer cannot replace the conventional computer, as used in digital photo and music.
- But, we should look further in n-qubit computers, with physical meanings and mathematical ideas.

Def. An *n-qubit* (= a vector state) is regarded as a 1-dimensional complex linear subspace of the dim 2^n Hilbert space, which can also be identified as a rank-1 projection in the form as a $2^n \times 2^n$ complex matrix.

Density Matrices

- In the formal setting of **non-commutative** probability, the **random** position of an n-qubit can be regarded as a density matrix, to be defined as a convex combination of rank-1 projections in M_2^n .

Def: A **density matrix** is a positive semidefinite matrix of trace 1.

- Thus, each 2 x 2 density matrix is expressed as

$$\frac{1}{2} \begin{bmatrix} 1-x & y+iz \\ y-iz & 1+x \end{bmatrix}$$

with $x^2 + y^2 + z^2 \leq 1$; so all density matrices fill up the whole solid sphere with S^2 as boundary.

- Nevertheless, for the case $n > 2$, there is no easy geometrical picture for the collection of all $n \times n$ density matrices.

Recap of the *simplest* Quantum Computer

- The setting of 1-qubit computer is a solid sphere in \mathbb{R}^3
--- just like the solid Earth in space.



- To send out quantum information ----- to *communicate* between two 1-qubit computers, we consider a feasible *affine transform* (preserving the 3-dimensional convex structure) of the solid **Earth**.

AFFINE TRANSFORMS induces LINEAR MAPS

- $M_n = M_n^+ - M_n^+ + i M_n^+ - i M_n^+$
- $M_n^+ = \mathbb{R}^+ \times \{\text{density matrices}\}$
- $\{\text{affine transforms on density matrices}\}$
 $\approx \{\text{trace-preserving positive linear maps}\}.$
- $\{\text{affine transforms on density matrices, fixing the scalar matrix}\}$
 $\approx \{\text{unital trace-preserving positive linear maps}\}.$
- $\{\text{feasible affine transforms on density matrices}\}$
 $\approx \{\text{trace-preserving completely positive linear maps}\}$
with deep unknown features of matrix analysis.

Shrew = Quantum Entanglements

of Positive Semi-Definite Matrices

Who's Afraid of

Quantum Entanglements?

Tensor - product setup for the **Taming** of the **Shrew**

- Consider a Hilbert space

$$H = H_1 \otimes H_2 .$$

- Some natural / simple / easy phenomena on H could be **entangled** in H_1 and H_2 separately.
- We wish to control the whole situation, bypassing / conquering /ignoring the entanglements.

Math Settings

- ❖ $L^2(X \times Y) = L^2(X) \otimes L^2(Y)$.
- ❖ Often consider of finite-dimensional Hilbert spaces as \mathbf{C}^n with a positive integer n .

➤ Thus $\mathbf{C}^n \otimes \mathbf{C}^k = \mathbf{C}^{nk}$.

$M_n =$ linear maps from \mathbf{C}^n to \mathbf{C}^n

$$M_n \otimes M_k = M_{nk} = M_n (M_k) = M_k (M_n).$$

--- no need to mention of anything as the **universal** property.

- In such an easy mathematical setting, who is afraid of **quantum entanglements** and **local-global** effects with respect to



Math Settings $M_n \otimes M_k = M_{nk}$ (with $n > 1, k > 1$)

➤ *{ the sums of $A_j \otimes B_j$ with A_j in M_n^+ , B_j in M_k^+ }
is only a proper subset of $(M_n \otimes M_k)^+ = M_{nk}^+$.*

Reason: $M_n^+ = \{ \text{positive linear combinations of rank-1 projections} \}$

- There are many rank-1 projections in M_{nk} which are not tensor product of rank-1 projections.

➤ Along this line, completely positive linear maps can go through the **quantum entanglements**, while positive linear maps cannot.

Quantum Entanglements provide **exciting** features for positive linear maps

Structure Theory

Notation: Each linear map $\varphi : M_n \rightarrow M_k$ can be extended to a linear map

$$\varphi \otimes \text{id}_p : M_n \otimes M_p \longrightarrow M_k \otimes M_p .$$

Def: φ is said to be *p-positive* when $\varphi \otimes \text{id}_p$ is a positive linear map.

Def: φ is said to be *completely positive* when φ is a *p*-positive linear map for each positive integer *p*.

Structure Theory

Thm (Choi) : All p -positive linear maps from M_n to M_k are **completely** positive when $n \leq p$ or $k \leq p$.

- Nevertheless, various p provide distinct classes of p -positive linear maps as elaborated in the following:

Example (Choi): The linear map $\varphi : M_n \rightarrow M_n$ defined as $\varphi(A) = (n-1)(\text{trace } A)I_n - A$ is $(n-1)$ -positive but not n -positive.

Main Thm: (Choi, 1975) A linear map

$\varphi : M_n \rightarrow M_k$ is ***completely positive***

iff $[\varphi(E_{ij})]_{i,j}$ is positive

where $\{E_{ij}\}$ are the matrix units

iff $\varphi(A) = \sum V_j^* A V_j$ for all $A \in M_n$

with $n \times k$ matrices V_j

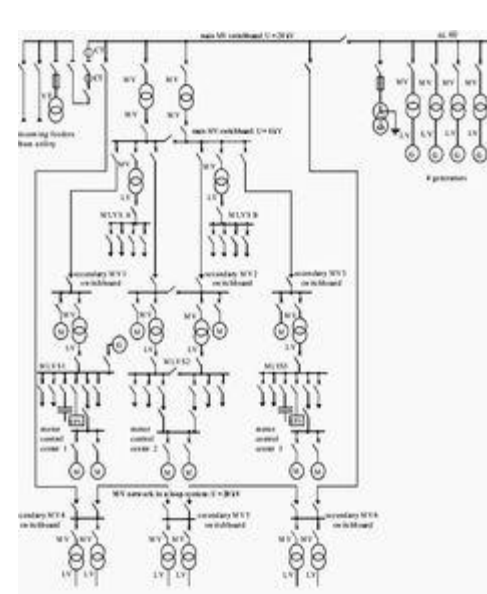
- This 1975 paper (6 pages) has been cited in nearly 3000 research papers, as of 2023 September

 Google Scholar

- More than 2000 citations in publications of Quantum Information.

CIRCUIT THEORY

- Each transformer defines a positive linear map $A \rightarrow V^*AV$.
Thus several transformers in series define a completely positive linear map.
- Main concern in circuit theory: General linear maps of mathematical expressions in terms of $[\varphi(E_{ij})]_{i,j}$ are not implementable.



➤ Classical computer vs Quantum computer

- ❖ A classical computers produces 0-1 sequences while a quantum computers produces psd matrices. Thus only completely positive maps are usable to connect *Quantum computers*

The Main Thm (Choi 1975) revisited

Let $\varphi : M_n \rightarrow M_k$ be a linear map. TFAE:

(1) φ is p -positive for all positive integer p .

(2) $[\varphi(E_{ij})]_{i,j}$ is positive

(3) $\varphi(A) = \sum V_j^* A V_j$ for all $A \in M_n$ with $n \times k$ matrices V_j

- (1) means to be the **hardest** nature to conquer all incredible quantum entanglements in $(M_n \otimes M_p)^+$ of various p .
- (2) is intended for the **simplest** mathematical expression of a general linear map.
- (3) turns to be the **only possible** connection in circuit theory.
- ❖ Stinespring Theorem (1955) covers the case (1) \Leftrightarrow (3).
- ❖ Theorem 1975 says much about (2) \Leftrightarrow (3) and (2) \Leftrightarrow (1), which is most needed in theory of quantum information.

Taming of Shrews

- NO way to describe so many incredible *entanglements* in $(M_n \otimes M_p)^+$ of various p .
- The most outstanding $T = \sum E_{ij} \otimes E_{ij} \in (M_n \otimes M_n)^+$, is a well behaved *entanglement* which serves as the representative for **ALL** wild entanglements.
- THEOREM says that to tame **ALL** shrews (= entanglements) is equivalent to tame a **single** LOVELY shrew (without worrying how nasty/dirty/undisciplined of other shrews).



of the **LOVELY** *Shrew*

Example $n=3$, $T = \sum E_{ij} \otimes E_{ij} \in (M_3 \otimes M_3)^+ = M_9^+$

❖ T is the **NATURAL** assemblage of matrix units

$$\begin{bmatrix} 1 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 1 \end{bmatrix}$$

➤ Indeed, $T^2 = nT$, so $\frac{1}{n}T$ is a rank-1 projection, but

T serves as the best witness to test all completely positive linear maps $M_3 \rightarrow M_3$.



Why Not Down to $n=2$?

- The **simplest** example of **quantum entanglement** is

$$\begin{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} & \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix} \\ \begin{bmatrix} 0 & 0 \\ 1 & 0 \end{bmatrix} & \begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix} \end{bmatrix}$$

as a positive 4 x 4 matrix, but not of the form as the sum of $A_j \otimes B_j$ with A_j in M_2^+ and B_j in M_2^+ .



Purpose: Wish to **classify** all linear maps

$\varphi : M_2 \rightarrow M_2$ by means of the 4 x 4 **Choi Matrix** C_φ

$$\begin{bmatrix} \varphi\left(\begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix}\right) & \varphi\left(\begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix}\right) \\ \varphi\left(\begin{bmatrix} 0 & 0 \\ 1 & 0 \end{bmatrix}\right) & \varphi\left(\begin{bmatrix} 0 & 0 \\ 0 & 1 \end{bmatrix}\right) \end{bmatrix}.$$

Challenge: *What sort of non-commutative geometry could be hidden/shown in the 4 x4 matrix C_φ ?*

Newest Classification Theorem

(Joint work with C.K. Li, 2023, JQIC)

Consider all $\varphi : M_2 \rightarrow M_2$ as unital trace-preserving and hermitian-preserving linear maps.

Then the 4 real eigenvalues of the Choi Matrix C_φ determine *the linear map φ up to unitary equivalence.*

I.e., iff C_φ and C_ψ have the same eigenvalues, then there exist unitaries U and W such that $\varphi(A) = U^ \psi(W^* A W) U$ for all A in M_2 .*

The Most Important Example:

By means of Pauli Matrices

$$Z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}, \quad X = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}, \quad Y = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix},$$

and 4 real numbers λ_j with $\sum \lambda_j = 1$.

Define $\varphi : M_2 \rightarrow M_2$

$$\text{as } \varphi(A) = \lambda_1 A + \lambda_2 ZAZ + \lambda_3 XAX + \lambda_4 YAY$$

- Then φ is a unital linear map preserving traces and hermitian matrices.
- The Choi Matrix $\mathbf{C}\varphi$ has $\{2\lambda_j\}$ as four eigenvalues.

Newest Classification Theorem

Restated

Each unital qubit channel φ

(unital trace preserving completely positive linear map $M_2 \rightarrow M_2$)

is unitarily equivalent to a **concrete** map of the

form $A \rightarrow \lambda_1 A + \lambda_2 ZAZ + \lambda_3 XAX + \lambda_4 YAY,$

where X, Y and Z are Pauli Matrices;

$\{2\lambda_j\}$ are eigenvalues of the Choi Matrix $C\varphi$.

➤ This provides the **WHOLE** picture of **unital qubit channels**.

OPEN QUESTION

What would be next *Classification Theorems* ?

Want to study the
case $n=3$.

- Need to understand the quantum entanglement of

1	0	0	0	1	0	0	0	1
0	0	0	0	0	0	0	0	0
0	0	0	0	0	0	0	0	0
0	0	0	0	0	0	0	0	0
1	0	0	0	1	0	0	0	1
0	0	0	0	0	0	0	0	0
0	0	0	0	0	0	0	0	0
0	0	0	0	0	0	0	0	0
1	0	0	0	1	0	0	0	1